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- Minimization of the energy with respect to all 267 backbone and side-chain dihedral angles is very time-consuming. We found it more efficient to minimize in two cycles. In the first, the energy was minimized with respect to all backbone dihe-
- dral angles and all side-chain dihedral angles  $\chi^1$  and  $\chi^2$  (which have the largest effect on the orientation of the side chain), i.e., 196 variables. In the second cycle, the energy was minimized with respect to all 155 side-chain dihedral angles. In one trial, 227 dihedral angles were varied (φ, ψ, χ¹, χ², and χ³), but in all other trials only 155 or 196 dihedral angles were varied.

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Differential Geometry and Polymer Conformation. 3. Single-Site and Nearest-Neighbor Distributions, and Nucleation of Protein Folding<sup>1</sup>

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ABSTRACT: A differential geometric representation of polymer chains is used to study protein backbone structure on the four- and five- $C^{\alpha}$  length scales. The analysis of the distribution of four- $C^{\alpha}$  units in the curvature-torsion  $(\kappa, \tau)$  plane reveals that there are islands of structure which occur with greater-than-random probability. These correspond to the well-known types of backbone structure—extended (E), right-handed  $\alpha$ -helical (A<sub>R</sub>), and flat-bend (A<sub>0</sub>) structure. It is shown that there are three distinct types of extended four- $\mathbb{C}^{\alpha}$ structure—left-handed twisted  $(E_L)$ , right-handed twisted  $(E_R)$ , and nearly flat  $(E_0)$ . The  $E_0$  and  $E_L$  structures form a structural continuum, but the  $E_R$  region is separated from this continuum by a region of low occupation. The A<sub>R</sub> and A<sub>0</sub> regions also form a continuum. These high-frequency islands are distributed throughout the occupied region of the  $(\kappa, \tau)$  plane and include structures of all types except that in the  $A_L$  (left-handed  $\alpha$ -helical) region. The high frequency of occurrence of these structures suggests that they are of low energy and therefore likely to occur in nucleation structures in the folding of the denatured molecule. It follows from the fact that these high-frequency structures occur throughout the occupied region, however, that there is relatively low selectivity in nucleation on the four- $C^{\alpha}$  scale, with structures representative of the entire occupied region being potential nuclei on the four-Ca scale. Extension of the analysis to the five-Ca scale shows that the high-frequency structures on this scale are made up of combinations of the high-frequency four- $C^{\alpha}$  structures and that a higher degree of selectivity in nucleation appears on the five-Ca length scale than on the four-Ca scale. Substantial differences are noted in the frequency of occurrence of the various combinations. Within the extended region, for example, the most frequently occurring structures are E<sub>0</sub>E<sub>0</sub>, with E<sub>R</sub>E<sub>R</sub> being the least frequent. A study of the correlation between nearest-neighbor four-Ca structures reveals that certain pairs of four- $C^{\alpha}$  structures have a low probability of occurring.  $A_0A_R$  and  $A_RA_0$  have a low positive correlation, while E<sub>X</sub>A<sub>R</sub> and A<sub>R</sub>E<sub>X</sub> (where X is R, L, or 0) have a negative correlation, indicating that these structures tend to avoid pairing. Five-C° components of nucleating structures are likely to be those which both occur with high frequency and show positive correlation between their four-Ca components. Of these, the most frequently occurring are ExEx and ARAR, which are repeating (regular) structures. The Ao four-Ca structure plays an important role in nucleation, despite its relative numerical infrequency, because it is the principal four- $C^{\alpha}$  structure which forms potential five- $C^{\alpha}$  nucleating structures which are nonregular. Larger nuclei, which are not considered explicitly in this paper, presumably arise as additional four-Ca units associated with five- $C^{\alpha}$  nuclei already present in the folding chain.

## I. Introduction

One of the significant aspects of protein architecture is the presence, in molecules with widely differing function and amino acid sequence, of certain well-defined structural features. Historically, the first such features to be noted were the  $\alpha$  helix and the extended strand (which associates with other strands to form the  $\beta$  sheet). These two structures are particularly easy to observe and classify because they are built up of repeating units (residues) with similar conformation. They are therefore capable of occurring, and of being observed, on a wide variety of backbone length scales, from a single residue to tens of residues.

It later became clear that there is at least one important structural feature which occurs on a single, well-defined length scale. This is the bend, or chain reversal, whose

presence is determined entirely by the structure of a backbone segment consisting of four successive a carbons—or, equivalently, by the values of two successive sets of  $(\phi, \psi)$  dihedral angles. Unlike  $\alpha$  helices or extended strands, bends are not built up of residues with essentially repeating values of  $(\phi, \psi)$ . In fact, a rather wide variety of combinations of  $(\phi_i, \psi_i; \phi_{i+1}, \psi_{i+1})$  form bends,<sup>3,4</sup> with quite similar  $C^{\alpha}$  arrangements.<sup>5</sup> Hence, bends differ from  $\alpha$  helices and extended structures (in  $\beta$  sheets) in the manner in which they are stabilized relative to their single-residue energies.  $\alpha$  helices and extended structures are each stabilized by (various interactions, including) a single type of interresidue hydrogen bond. Bends, on the other hand, may be stabilized by hydrogen bonds in a number of ways<sup>6</sup> or involve no hydrogen bonds at all, depending on the particular values of  $(\phi, \psi)$  which obtain. It is also possible for bends to be stabilized in still other ways—e.g., by hydrophobic interactions—completely external to the bend region itself.<sup>7</sup>

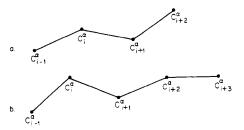
One then can raise the following question: Are there other well-defined structural features in proteins, as yet unobserved because they occur on length scales which are not easily accessible through conventional structure representations?

The identification and characterization of specific structural features of proteins is important because such structures are believed to be essential for the correct self-organization of the protein molecule, a process which is crucial to its proper function. For example, it has been proposed<sup>8</sup> that the nucleation regions which initiate folding are hydrophobic pockets that form around bends; another theory of folding<sup>9</sup> suggests that nucleation occurs by formation and disintegration of  $\alpha$  helices.

Before proceeding, it is necessary to define the term nucleation, as we shall use it in this paper. The denatured protein may be regarded as a fluctuating "random coil". We suggest that, as renaturing conditions are created, ordered structures form in the fluctuating chain; their lifetimes are much longer than those of any conformations which form in neighboring parts of the chain. We regard these long-lived structures as nuclei. These structures need not all form simultaneously in the fluctuating coil, and, indeed, it may be necessary for them to form in a specific sequence to attain the correct globular folding. Furthermore, not all of these structures need individually cause a reduction in chain dimensions. Previous work<sup>8</sup> has been directed toward the identification of the most stable structure responsible for a decrease of chain dimensions. We shall be interested in the entire spectrum of possible nucleating structures.

Nucleating structures of the type that we have described can be of various lengths. Such structures must be built up residue-by-residue, beginning from very short lengths of backbone which assume the proper conformation solely under the influence of localized, short-range interactions. Thus, each nucleus undergoes an elongation process until it reaches its ultimate length. It then contributes to the folding process either by bringing about favorable interactions between other parts of the chain or by participating in such interactions itself.

In this paper, we use a differential geometric (DG) representation of polypeptide chains to extend the search for significant structures systematically to longer length scales than that treated by the  $(\phi, \psi)$  representation. As was demonstrated in papers  $1^{10}$  and  $2^5$  of this series, the DG representation operates on a four- $C^{\alpha}$  length scale (Figure 1a). It is therefore able, for example, to give a structurally clear description of bends, which is not easily



**Figure 1.** (a) Four- $C^{\alpha}$  unit whose conformation is described by  $(\kappa_i, \tau_i)$ . (b) Five- $C^{\alpha}$  unit whose conformation is described by  $(\kappa_i, \tau_i; \kappa_{i+1}, \tau_{i+1})$ . These are representations of virtual-bond structures, and the quantities  $\kappa_i$  and  $\tau_i$  were defined in the first paper. <sup>10</sup>

provided by the  $(\phi, \psi)$  representation.<sup>5</sup> We shall begin by investigating the distribution of values of curvature and torsion  $(\kappa_i, \tau_i)$  of four- $C^{\alpha}$  units (defined in paper  $1^{10}$ ), with a view to characterizing structure on this scale in a systematic fashion. We shall then investigate the distribution of nearest-neighbor  $(\kappa_i, \tau_i, \kappa_{i+1}, \tau_{i+1})$  pairs. This distribution provides information about the structure of five- $C^{\alpha}$  units (Figure 1b) and represents a first step in the extension of the characterization of significant structures to still longer length scales.

We shall adopt the hypothesis that those structures that occur with high frequency in native proteins are of low energy and therefore have a high probability of forming in the denatured chain under renaturing conditions. We shall identify these as possible nucleation structures. This is not to suggest that four- or five- $C^{\alpha}$  structures are the complete nucleating structures. This is possibly, but not necessarily, true. Examination of protein structures on a variety of length scales should reveal that some high-frequency structures have well-defined length scales and others do not-precisely the situation that we noted above in connection with bends, helices, and extended strands. We do suggest, however, that the four- and five- $C^{\alpha}$  nucleating structures that we shall identify form parts of complete nucleating structures and therefore are potentially important in the process of growth of nuclei. Their formation constitutes an early stage in the growth of the individual nuclei. Thus, our goals in this paper are (1) to delineate the structural types which are significant on the four- $C^{\alpha}$  and five- $C^{\alpha}$  length scales: (2) to provide a quantitative picture of the relative importance of these different structural types in native proteins; and (3) to draw conclusions from this analysis about the nucleation process which leads to correct folding of the protein.

#### II. Single-Site Distribution

In paper 1,<sup>10</sup> we explored the localization of various ordered backbone structures in  $(\kappa, \tau)$  space by determining  $(\kappa_i, \tau_i)$  at residues which had been determined by  $(\phi, \psi)$  criteria to be in the conformational state of interest. We now adopt a more general approach and ask what types of characteristic structures occur on the single-site (four- $C^{\alpha}$ ) length scale in  $(\kappa, \tau)$  space, without reference to the values of  $(\phi, \psi)$ . [We shall use the term "single site" throughout this paper to denote an individual  $(\kappa_i, \tau_i)$  pair, ascribed by convention<sup>10</sup> to residue *i*. It should be borne in mind, however, that  $(\kappa_i, \tau_i)$  is determined by the locations of the four atoms  $C^{\alpha}_{i-1}$  to  $C^{\alpha}_{i+2}$  (Figure 1a) or alternatively by two pairs of dihedral angles  $(\phi_i, \psi_i; \phi_{i+1}, \psi_{i+1})$ .]

This more general approach is particularly important because of the degeneracy in terms of the values of  $(\phi, \psi)$  which is observed on the four- $C^{\alpha}$  length scale.<sup>5</sup> This degeneracy implies that very similar four- $C^{\alpha}$  virtual-bond structures, having widely differing combinations of  $(\phi_i, \psi_i; \phi_{i+1}, \psi_{i+1})$ , can exist. It follows that an analysis in terms

- 1. bovine pancreatic trypsin inhibitor
- concanavalin A
- 3. carboxypeptidase A
- 4. clostridial flavodoxin
- 5. thermolysin
- 6. oxidized chromatium high-potential iron protein
- 7. staphylococcal nuclease
- 8. sea lamprey hemoglobin
- 9. sperm whale myoglobin
- 10. hen egg white lysozyme
- 11. ribonuclease S
- 12. subtilisin BPN' (Novo)
- 13. rubredoxin
- 14. carbonic anhydrase
- 15. tuna cytochrome c (oxidized inner)
- 16. D-glyceraldehyde 3-phosphate dehydrogenase
- 17. papain
- 18. carp myogen
- 19. human plasma prealbumin
- 20. ferrodoxin
- 21.  $\alpha$ -chymotrypsin
- chicken triose phosphate isomerase

of  $(\phi, \psi)$  may well fail to reveal certain structural types which, though well-defined on the  $(\kappa, \tau)$  length scale, are diffuse in terms of the values of  $(\phi, \psi)$ . An obvious example of such a structural type is the bend, which was discussed in the Introduction. The traditional  $(\phi, \psi)$  classification, <sup>3,4</sup> which produces 11 bend types, does not easily reveal the fact that bends form a continuum of structural types, ranging from segments of  $\alpha_R$ -helical structures through flat bends to segments of  $\alpha_L$ -helical structures.<sup>5</sup> There is every reason to suppose that the same effect holds in other regions of  $(\kappa, \tau)$  space.

We adopt a graphical approach to the problem and begin by constructing the observed distribution of values of  $(\kappa, \tau)$  in a sample of proteins from their X-ray coordinates. A preliminary form of this distribution was exhibited in Figure 3 of paper 1. To For the current study, the sample of 13 proteins used in paper 1 was expanded by the addition of nine more proteins, so chosen that, in the combined sample (as in the original sample of paper 1), there are few homologies and the coordinates available are of reasonable quality. The proteins used are listed in Table I; all coordinates were obtained from the Brookhaven Protein Data Bank.

The resulting distribution is shown in Figure 2, plotted as the number of four- $C^{\alpha}$  segments per  $(0.1 \times 0.1)$  rad/Å square in the  $(\kappa, \tau)$  plane, independent of residue type. This is essentially a top view of a three-dimensional bar graph, in which the heights of the bars are given by the occupancy numbers of the squares. The essential identity of the occupied region of Figure 2 with that of the smaller sample is readily seen in Figure 2, in which the occupied squares of the graph have been assigned an arbitrary numbering for convenient reference (small numbers in lower right hand corners of squares). Squares 1-66 are occupied in the original 13-protein sample of paper 1. Squares 70-73 and 76 are occupied only in the 22-protein sample of the current work. (The omission of some numbers from the arbitrary numbering of squares arose from an earlier version of this work and is of no significance.) There are only seven points in the five new squares. Comparison of the two distributions shows that there is also no significant change in the relative proportion of points falling in the various regions of the  $(\kappa, \tau)$  plane, in passing from the original sample of 2127 points to the current sample of 3754 points. We therefore believe that the two samples both represent reasonably well the actual

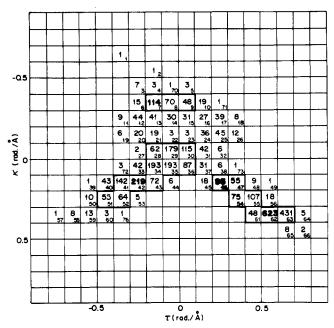


Figure 2. Untransformed single-site distribution of values of  $(\kappa,\tau)$  for the sample of 22 proteins listed in Table I. The islands of greater-than-random occupation are indicated by heavy outlines. The large number in each  $(0.1\times0.1)$  rad/Å square is the number of points observed to fall in that square, and the small number is an arbitrarily assigned numbering of the squares. Occupation numbers in boldface indicate that those squares are maxima in the distribution. Shaded area (squares 46, 47, and 54) indicates one of the two peaks on an island with two maxima (see text and Appendix II). The squares comprising the various regions of peaks in this distribution  $(A_0,\,E_L,\,$  etc.) are identified in Table II.

distribution of conformational preferences in proteins. Before proceeding, it is necessary to recall briefly a number of facts about the  $(\kappa, \tau)$  plane which were developed in paper  $2^5$  (to which the reader is referred for mathematical details). Two in particular will be of interest.

1. There is a discontinuity in  $\kappa$  and  $\tau$  as the conformation of a four-C<sup>α</sup> segment is changed continuously, at a conformation where the factor  $p_i p_{i+1}$  (which enters into the formulas for  $\kappa_i$  and  $\tau_i$ ) changes sign  $(p_i p_{i+1})$  assumes the values  $\pm 1$ ). The quantity  $p_i$  is defined in eq 6 of paper  $1^{10}$ as the scalar product of two nonperpendicular unit vectors at the ith and (i-1)th  $\alpha$ -carbon atoms, respectively; it is a factor which guarantees that the unit normal in the local reference frame at residue i has a positive projection on the unit normal in the local reference frame at residue i -1. As a result, there are certain regions of the  $(\kappa, \tau)$  plane where structures that are conformationally very similar appear as widely separated points. In paper 2, it was shown that the structures affected by this discontinuity are bends and allied structures. It was also shown that a transformation exists (eq 16 of paper 2) which brings structures for which  $p_i p_{i+1} = +1$  (which are in the minority) to new points in the  $(\kappa, \tau)$  plane which are adjacent to those representing conformationally similar structures for which  $p_i p_{i+1} = -1$ . In Appendix I, we apply this transformation to the distribution of Figure 2 and perform an analysis of the resulting distribution which demonstrates that the conclusions that we shall draw from Figure 2 are correct. This is important because we shall use the untransformed values of  $(\kappa, \tau)$  to construct the nearest-neighbor distribution of section III. We do this because the untransformed distribution provides a natural method for distinguishing between the right- and left-handed twisted (AR and A<sub>L</sub>) bends, on the one hand, and the structurally adjacent flat (A<sub>0</sub>) bends, on the other. We shall show that these structures have distinctly different behavior in the nearest-neighbor distribution.

2. There is also a discontinuity in  $\kappa$  for four-C° structures with  $p_i p_{i+1} = -1$  at the conformational point where  $\kappa$  changes sign. It was demonstrated in paper 2 that this leads to a butterfly-shaped gap centered about the  $\tau$  axis.

Both of these factors will be taken into account (although not explicitly) in the analysis of the distributions.

We now return to the analysis of Figure 2. If the sample of 3754 points were distributed uniformly over the 71 squares which constitute the occupied region of Figure 2. as might be if each four-Ca unit were equally likely to assume any conformation in the occupied region, the population per square would be 52.9.12 We therefore ask which squares are occupied with greater-than-random frequency. These squares form three separate islands. which are indicated in Figure 2 by heavy outlines. [To account for possible statistical error, we included squares with population  $\geq 48$  (=0.9 × 53), where the factor 0.9 was chosen arbitrarily. It should be noted that, whereas the squares included in these islands contain only 30% of the available area, they contain 81.1% of the residues indicating the utility of the island approach for identifying high-frequency structures.

What types of four- $C^{\alpha}$  structure are represented by the points contained in the islands of high frequency? The answer to this question becomes clear from a comparison of Figure 2 with Figures 6–8 of paper 1. The island embracing squares 7–9 (which we shall refer to as island I) corresponds precisely to the low-torsion bend region of Figure 8, paper 1. The island which includes the origin (which we shall refer to as island II) substantially overlaps the distribution of extended structures of Figure 7, paper 1. The remaining island (island III) includes the principal  $\alpha_{\rm R}$ -helical region of Figure 6, paper 1, and the (coincident)  $\alpha_{\rm R}$ -helical bend region of Figure 8. It also includes part of the region identified as including extended structures in Figure 7, paper 1.

A more detailed analysis of the individual islands is in order. We begin with island II, which includes the largest part of the extended region. There is a striking breadth of structural types included in this island. These range from structures which are very nearly planar zigzags (such as those concentrated near the origin, in squares 29, 30, 35, and 36) to those which are nearly  $\alpha_{\rm L}$ -helical (square 51). There is a very strong predominance of left-handed structures. The population maximum at square 42 falls in a region which is both bent (in terms of change of chain direction, as shown by  $\kappa$ ) and twisted (in a left-handed sense, as shown by  $\tau$ ) relative to a perfect planar zigzag form.

There is only one maximum in island II (that at square 42), and this island clearly contains a continuum of structural types. Nevertheless, it is useful and meaningful to distinguish between the nearly planar structures centered about the origin and the more twisted, more bent structures clustered about the maximum. We shall denote these as  $E_0$  and  $E_L$  structures, respectively. The subscript in  $E_0$  is intended to denote the fact that these structures are almost flat; of course, since they are usually not perfectly planar, they do have a handedness. The  $E_L$  structures have a pronounced left-handed twist. <sup>13-19</sup> We use a capital E to emphasize the fact that the structures defined by these values of  $(\kappa, \tau)$  are on a longer scale than those defined by a  $(\phi, \psi)$  pair which falls in the extended ( $\epsilon$ ) region of the ( $\phi$ ,  $\psi$ ) plane. To make our notation precise, we define as E<sub>0</sub> those structures which are represented by points which fall in squares 28-30 and 34-36. E<sub>L</sub> structures

are those which fall into squares 41–43 and 52. Because the main body of the  $A_L$  region (squares 58–60) is sparsely occupied, square 51 is assigned to the  $E_L$  region, although the symmetric right-handed square (55) is assigned as an outlier of the  $A_R$  region. This flexibility of assignment illustrates strikingly the general fact that there is a continuous gradation of four- $C^\alpha$  structural types, rather than a collection of discrete varieties.

An obvious question is whether there is a corresponding E<sub>R</sub> structure in the region of high-frequency structures. One would expect such structures to have the same values of  $\kappa$  as  $E_L$  structures but values of  $\tau$  opposite in sign. Inspection reveals that the maximum at square 46 in island III and the neighboring squares of that island (i.e., squares 47 and 54) indeed constitute the E<sub>R</sub> region. It should be noted that the populations of squares 46 and 47 are very significantly lower than that of the corresponding E<sub>L</sub> region, but those of squares 54 and 55 are greater than the population of the corresponding left-handed region. The curvature of squares 54 and 55 is greater than that at squares 46 and 47, as is the torsion, and these structures are substantially closer to the A<sub>R</sub> region (we use capital A rather than  $\alpha$  for the same reason that we used capital E above). Thus the gradual change in parity (handedness) preference from left- to right-handedness as we pass continuously along the positive  $\kappa$  axis from the E region, near the origin, to the A (helical) region is clearly manifest; i.e., near the origin, left-handedness predominates but, as we move away from the origin in the direction of increasing κ, right-handedness becomes increasingly preferred.

It should also be noted that island III is separated from island II by a channel of less-than-random occurrence (squares 31, 37, and 45) which falls in the extreme right-handed section of the  $E_0$  region (opposite in handedness to squares 28, 34, and 43). This emphasizes again the strong preference for left-handedness<sup>13–19</sup> in the E region.

As was noted above, island III contains two local maxima, those at squares 46 and 62. The latter is the A<sub>R</sub>-helical peak, which, as might be expected, overwhelms in magnitude any other feature of the distribution. The maximum at square 62 has associated with it squares 55, 61, and 63; these together are defined as the A<sub>R</sub> region. (A general procedure for dividing islands containing more than one maximum into separate peaks is outlined in Appendix II.)

Because of discontinuity no. 1 in the  $(\kappa, \tau)$  plane, noted above, island I is actually contiguous structurally with the  $A_R$  region of island III (see also Appendix I) in the sense that four- $C^{\alpha}$  structures represented by points which fall in square 7 are very similar to those represented by points which fall in square 63 (see paper 2 for details). It is thus clear that a continuum of bend structures exists, as was shown in paper 2. (We denote the bend structures that occupy island I, which are more planar than those of the  $A_R$  region and include essentially flat bends near  $\tau=0$ , as  $A_0$  bends, in the same spirit as the  $E_0$  notation above.) In fact, island I is contiguous also with the  $A_L$  region (squares 9 and 57); however, the  $A_L$  region is less-than-randomly occupied (except for the borderline square 51, which was discussed above).

The outcome of this analysis is a remarkable result. The region of greater-than-random occupation forms a continuum (broken by only one gap) which includes four- $C^{\alpha}$  structures representative of the entire occupied region (and therefore presumably the entire accessible conformational region) except for the left-handed helical region. This suggests that nucleation sites can be assembled out of four- $C^{\alpha}$  "pieces" representing almost the full range of

Table II Characteristics of Peaks in Single-Site Distribution a

peak type	coordinates of maximum	contracted numbering of maximum	contracted numbering of squares in peak	population	đ	$ar{s}$
$A_0$	(6, 8)	7	7-9	232	0.72	1.65
$\mathbf{E_L^0} + \mathbf{E_0}$	(11, 7)	42	$(41-43, 51, 52)^b$ $(28-30, 34-36)^c$	1381 (552, <sup>b</sup> 829 <sup>c</sup> )	1.79	2.89
$\mathbf{E}_{\mathbf{R}}$	(11, 12)	46	46, 47, 54	223	0.72	1.47
$A_R$	(13, 15)	62	55, 61-63	1209	0.52	9.19

<sup>a</sup> All data pertain to Figure 2. <sup>b</sup>  $E_L$  region. <sup>c</sup>  $E_0$  region.

possible structures. There are, of course, significant differences in population between different types of structure within the region of greater-than-random occurrence. This provides an important measure of discrimination between different types of nucleation sites in the sense that certain structural types, which are observed as maxima (e.g., E<sub>L</sub>, E<sub>R</sub>, and A<sub>R</sub>), can be regarded as prototypical nucleation structures on the four- $C^{\alpha}$  length scale. Nevertheless, in view of the almost continuous distribution of frequently occurring structural types, it is clear that selectivity of nucleation is rather weak on the four- $C^{\alpha}$  length scale. It will therefore be of great interest to see what restrictions on pairing of four- $C^{\alpha}$  structures operate in the formation of structure on the five- $C^{\alpha}$  length scale. This problem will be taken up in the next section.

It should be noted that this result is independent of the size of the grid. An analysis of the distribution shown in Figure 2 using a grid twice as fine  $[(0.05 \times 0.05) \text{ rad/Å}]$ gives an essentially identical high-frequency region.

It is desirable to have a quantitative method for the representation of various properties of the peaks observed in bar graphs. For this purpose, we define the average value of a parameter D (which has the value  $D_i$  in square i) for a given peak as

$$\bar{D} = \sum_{j} D_{j} n_{j} / \sum_{j} n_{j} \tag{1}$$

where the sum runs over all squares which are included in the peak and  $n_i$  is the population of the jth square.<sup>20</sup>

We shall use two parameters to characterize the peaks appearing in Figure 2. They are d, the average distance of squares from the peak maximum, and  $\bar{s}$ , the average degree of nonrandomness of the peak. Each square in the  $18 \times 18$  grid of Figure 2 can be characterized by its row and column position. Thus, square 3 has coordinates (5, 7) and square 61 has coordinates (13, 14). The distance  $d_{12}$  between two squares, say  $(m_1, n_1)$  and  $(m_2, n_2)$  is given by the usual Cartesian formula:

$$d_{12} = [(m_1 - m_2)^2 + (n_1 - n_2)^2]^{1/2}$$
 (2)

where  $d_{12}$  is measured in units of 0.1 rad/Å, the "lattice constant" of the graph. If one of these squares is the peak maximum, and the appropriate average is taken as in eq 1,  $\bar{d}$  is obtained.<sup>21</sup> Clearly  $\bar{d}$  is a measure of the width of the peak. Small values of  $\bar{d}$  correspond to narrow peaks, and large values to broad peaks. It should be noted that  $s_i$ , the degree of nonrandomness of square i, is given by

$$s_i = n_i / 52.9 \tag{3}$$

and that  $s_i > 1$  corresponds to greater-than-random occurrence and  $s_i < 1$  to less-than-random occurrence. Substitution of  $s_i$  of eq 3 for  $D_i$  in eq 1 yields  $\bar{s}$ .

The parameters calculated for the peaks observed in the three islands of Figure 2 are shown in Table II. The E. +  $E_0$  peak (island II) is clearly the broadest, as indicated by its large value of d, 1.79. The  $A_0$  and  $E_R$  peaks have essentially the same width, with  $\bar{d} = 0.72$ . The narrowest

peak is the  $A_R$  peak, with  $\bar{d} = 0.52$ . (We stress again the fact that the AR and Ao peaks are structurally contiguous and that our separation of them is convenient, but somewhat arbitrary on the four- $C^{\alpha}$  length scale. In section III, we shall see that, on the five- $C^{\alpha}$  length scale,  $A_R$  and  $A_0$ structures exhibit rather different behavior.) These values of d reflect accurately the fact that the  $\alpha_R$  helix is the most sharply defined ordered backbone feature and the wide latitude which obtains in extended structures. The AR peak has  $\bar{s} = 9.19$ , reflecting the fact that the  $A_R$  structure, by virtue of its occurrence either alone, as a bend, or in a combination with other  $A_R$  structures as  $\alpha_R$  helices, <sup>22-24</sup> is overwhelmingly the most likely single, sharply defined structure to occur on the four- $C^{\alpha}$  length scale; it is followed in descending order of  $\bar{s}$  by the  $E_L + E_0$ ,  $A_0$ , and  $E_R$  peaks.

We can summarize the conclusions that we have drawn from the single-site (four- $C^{\alpha}$  structure) distribution as follows:

- 1. The structures which occur with greater-than-random frequency on this length scale form a virtual continuum. with only one break (in the channel formed by squares 31, 37, and 45 in the right-handed extended region), which includes structures representative of the entire occupied region of the  $(\kappa, \tau)$  plane, except for the left-handed helical region.
- 2. We postulate that all of these high-frequency structures on the four- $C^{\alpha}$  scale are available to act as components of nucleation regions in the folding of the protein. It therefore seems that there is weak selection of nucleation structures on the four- $C^{\alpha}$  length scale and that stronger selection must take place on longer length scales in the protein.8,25,26
- 3. The extended region is seen to include four- $C^{\alpha}$ structures with either right- or left-handed twists, as well as many structures which are almost planar. These are designated as E<sub>R</sub>, E<sub>L</sub>, and E<sub>0</sub> structures, respectively. The preference for left-handedness in the extended region is clearly demonstrated. 13-19
- 4. Parameters were developed to describe quantitatively the localization and population of the various peaks observed in the single-site (four- $C^{\alpha}$ ) distribution. It was shown that the A<sub>R</sub> peak is very narrow and very high, that the  $E_0 + E_L$  peak is quite broad and rather high, and that there are two lower peaks which are narrow—the A<sub>0</sub> (flat bend) and  $E_R$  peaks.

We now extend the methods that we have developed to the study of the nearest-neighbor pair distribution of values of  $(\kappa, \tau)$ , which reflects the structure of the backbone on the five- $C^{\alpha}$  length scale.

#### III. Nearest-Neighbor Distribution

The protein sample given in Table I leads to 3729 nearest-neighbor pairs<sup>27</sup> of values of  $(\kappa_i, \tau_i; \kappa_{i+1}, \tau_{i+1})$ , the distribution of which can be represented on a graph in which are given the populations of  $(0.1 \times 0.1 \times 0.1 \times 0.1)$ rad/Å hypercubes in the  $(\kappa_i, \tau_i; \kappa_{i+1}, \tau_{i+1})$  hyperplane. These four values, plus the height, constitute essentially a five-dimensional bar graph. It should be noted that it is also possible to represent the data in a three-dimensional contracted graph by assigning to each hypercube two coordinates corresponding to the arbitrary successive numbering that we have assigned to the occupied squares in the  $(\kappa, \tau)$  plane (Figure 2). We shall find this to be a convenient shorthand notation for identifying hypercubes. The contraction, however, eliminates the structural information contained in the five-dimensional representation, so that we cannot use it for structural studies; these must be carried out in the full five-dimensional space.

We shall find it convenient to have a notation which enables us to distinguish succinctly between the  $(\kappa, \tau)$  and  $(\kappa_i, \tau_i; \kappa_{i+1}, \tau_{i+1})$  planes and between these and the contracted representation of the latter. We will therefore refer to these henceforth as  $\mathbf{P}^{(2)}$ ,  $\mathbf{P}^{(4)}$ , and  $\mathbf{C}^{(2)}$ , respectively. The single-site (three-dimensional) and nearest-neighbor (five-dimensional) graphs will be denoted by  $\mathbf{G}^{(3)}$  and  $\mathbf{G}^{(5)}$ , respectively.

The fact that we are forced to operate in a five-dimensional space deprives us of the advantages afforded by geometrical intuition. We must therefore fall back on analytical methods which, in the three-dimensional case, reduce to intuitively evident procedures.

The first step in our analysis is to locate the maxima in  $G^{(5)}$ . This is done by searching for hypercubes in  $P^{(4)}$  whose population is greater than that of any of their nearest neighbors. (We note in passing that the number of nearest neighbors in n dimensions is given by  $3^n - 1$ , so that, whereas a given cube in  $P^{(2)}$  has 8 nearest neighbors, a hypercube in P<sup>(4)</sup> has 80.) In order to decide which of these peaks are of sufficient height to be significant, we observe that the 3729 pairs are distributed among only 1052 of the  $5041 (=71^2)$  theoretically accessible hypercubes, where 71 is the number of occupied squares in the single-site distribution of Figure 2. If the pairs were distributed randomly in this occupied subspace, there would be a uniform population per hypercube of 3729/1052, or  $3.54 (\approx 4)$ . We shall therefore accept as significant structural features those peaks with height  $h \ge 5$  and assign hypercubes to them by using an exactly analogous procedure (Appendix II) to that outlined in section II, with the proviso that the population  $n_i$  of such hypercubes be greater than 4 (to distinguish it from random occupancy).12 This procedure for constructing peaks may err on the side of conservatism in neglecting hypercubes with  $n_i \le 4$ ; nevertheless, it seems to be a reasonably satisfactory method for picking peaks out of the background.

Another factor that must be considered is that of correlation. It should be noted that, in constructing the nearest-neighbor graph  $G^{(5)}$ , we use a sample which also has the single-site distribution given by  $G^{(3)}$  (Figure 2). Thus, the degree of randomness of pairing, or degree of correlation, is given by

$$c_{ij} = s_{ij}/s_i s_j \tag{4}$$

where  $s_k = n_k/3754$  is the fraction of the single-site distribution occurring in square k and  $s_{kl} = n_{kl}/3729$  is the fraction of nearest-neighbor pairs occurring in hypercube kl (where we use the contracted numbering of Figure 2).  $c_{ij} > 1$  indicates that the nearest-neighbor pair ij occurs in  $\mathbf{G}^{(5)}$  with greater frequency than would be expected on the basis of the frequencies of i and j in  $\mathbf{G}^{(3)}$ ;  $c_{ij} < 1$  indicates the reverse.

As in the case of  $G^{(3)}$ , we shall be interested in the distance d of a given hypercube in  $P^{(4)}$  from the maximum with which it is associated. A new parameter of interest in the analysis of  $G^{(5)}$  is the distance  $\Delta_{ij}$  in  $P^{(2)}$  between

Table III

Peaks in G<sup>(5)</sup> Organized by Islands

island	coordinates of peak <sup>a</sup>	contracted coordinates	height <i>h</i>
I	(6, 8, 6, 8)	(7, 7)	8
II	(6, 8, 13, 15)	(7, 62)	19
	(6, 9, 11, 12)	(8, 46)	7
III	(7, 8, 9, 9)	(13, 29)	6
IV	(7, 9, 11, 7)	(14, 42)	6 8
	(8, 12, 11, 7)	(25, 42)	9
	(9, 9, 11, 6)	(29, 41)	20
	(10, 8, 9, 9)	(34, 29)	31
	(11, 6, 11, 7)	(41, 42)	16
V	(10, 9, 7, 8)	(35, 13)	7
VI	(10, 9, 13, 15)	(35, 62)	22
	(11, 7, 13, 15)	(42, 62)	25
	(11, 12, 13, 16)	(46, 63)	8
	(13, 15, 13, 16)	(62, 63)	261
VII	(12, 13, 11, 12)	(54, 46)	6
VIII	(11, 7, 11, 12)	(42, 46)	8
IX	(11, 12, 11, 7)	(46, 42)	9
	(13, 16, 12, 6)	(63, 52)	9
$\mathbf{X}$	(13, 15, 6, 8)	(62, 7)	28
XI	(13, 15, 8, 12)	(62, 25)	7
XII	(12, 13, 9, 8)	(54, 28)	8 <b>b</b>
	(12, 14, 10, 9)	(55, 35)	8 b

<sup>a</sup> Row and column numbering from Figure 2. <sup>b</sup> Double peak (see text).

the two squares which make up the nearest-neighbor pair. The smaller this distance is, the more nearly the hypercube in question represents a (regular) repeating structure. In analogy with eq 1, these parameters can all be used in the form of population-weighted averages over the hypercubes in a given peak:

$$\bar{a} = \sum_{ij} n_{ij} a_{ij} / \sum_{ij} n_{ij}$$
 (5)

In using these parameters, account must be taken of the mathematical features of the  $(\kappa, \tau)$  representation summarized in the previous section. This is particularly true because, as in the discussion of  $G^{(3)}$  in section II, and for the reason noted there, we will deal with the untransformed nearest-neighbor distribution characterized by mixed values of  $p_i p_{i+1}$  (=±1). [In Appendix III we demonstrate that the transformed  $(p_i p_{i+1} = -1 \text{ only})$  nearestneighbor distribution gives results which are completely consistent with the analysis in the main text.]

Of the 32 maxima found in  $G^{(5)}$ , there are 20 with  $h \ge 5$ . These fall on 11 islands composed of hypercubes with population greater than 4; the peaks are shown, organized by island, in Table III. There is one additional small island (no. XII of Table III) with a double peak, i.e., two adjacent hypercubes with the same population. As in the discussion of  $G^{(3)}$ , we remark that the total number of hypercubes contained in these islands is 164, or 15.6% of the total number of occupied hypercubes, whereas the total population of these islands is 2219 nearest-neighbor pairs, or 59.5% of the total number of pairs. This indicates clearly that the usefulness of the island approach for the determination of significant structural features is retained in  $G^{(5)}$ .

The peaks occurring in  $\mathbf{G}^{(5)}$  can be classified according to the types of structure they represent, using as a basis the classification of four- $\mathbf{C}^{\alpha}$  units established in section II. The results are shown in Table IV.

We begin our discussion of Table IV by considering the EE peaks, which represent five- $C^{\alpha}$  structures which fall entirely in the extended region. It can be seen that nearly all possible combinations of the three extended types that we defined previously ( $E_R$ ,  $E_L$ , and  $E_0$ ) occur. It is instructive to examine the relative importance of the dif-

Table IV Peaks in G(5) Organized by Type

	peak type	coordinates of maximum	contracted coordinates of maximum	$N_{\mathbf{c}}{}^{a}$	h <sup>b</sup>	$\overline{c}$	$\overline{d}$	$\overline{\Delta}$	population of peak
	,EE	(8, 12, 11, 7)	(25, 42)	8	9	3.55	1.31	4.77	51
		(11, 12, 11, 7)	(46, 42)	6	9	2.43	1.22	6.22	41
	$\begin{array}{c} \mathbf{E_R} \mathbf{E_L} \\ \mathbf{E_L} \mathbf{E_R} \end{array}$	(11, 7, 11, 12)	(42, 46)	2	8	2.12	0.54	5.41	13
	E <sub>LE</sub> L	(11, 6, 11, 7)	(41, 42)	16	16	2.03	1.43	1.57	128
EE	$\mathbf{E} \left\langle \mathbf{E}_{\mathbf{R}}^{\mathbf{L}} \mathbf{E}_{\mathbf{R}}^{\mathbf{L}} \right\rangle$	(12, 13, 11, 12)	(54, 46)	1	6	3.26	0.00	1.41	6
	$\mathbf{E}_{0}\mathbf{E}_{L}$	(9, 9, 11, 6)	(29, 41)	17	20	2.39	1.29	2.93	163
	$\mathbf{E}_{0}^{\mathbf{E}}\mathbf{E}_{0}^{\mathbf{E}} + \mathbf{E}_{\mathbf{L}}\mathbf{E}_{0}^{\mathbf{c}}$	(10, 8, 9, 9)	(34, 29)	37	31	2.31	1.60	1.65	444
						$(2.40, 2.04)^c$	(1.46, -)	(1.09, 3.24)	(328, 116)
	$^{L}_{R}\mathrm{E}_{o}$	(12, 13, 9, 8)	(54, 28)	2	8	4.40	0.77	4.80	13
	$\int_{\mathbf{E}_0} \mathbf{A_R}$	(10, 9, 13, 15)	(35, 62)	8	22	0.98	0.81	6.34	63.5
EAD	ETAR	(11, 7, 13, 15)	(42, 62)	9	25	0.80	1.00	8.83	81.5
1	$\mathbf{E}_{\mathbf{E}_{\mathbf{R}}\mathbf{A}_{\mathbf{R}}}^{\mathbf{E}_{\mathbf{L}}\mathbf{A}_{\mathbf{R}}}$	(11, 12, 13, 16)	(46, 63)	5	8	0.84	1.03	3.66	32
	$A_R E_0$	(12, 14, 10, 9)	(55, 35)	1	8	1.45	0.00	5.39	8
$A_R E$	EXA <sub>D</sub> E <sub>T</sub>	(13, 16, 12, 6)	(63, 52)	4	9	0.69	0.80	9.23	<b>2</b> 8
	$A_R^R E_R$	(13, 15, 8, 12)	$(62, 25)^d$	1	7	0.95	0.00	5.83	7
	$\mathbf{E}_{o}\mathbf{A}_{o}$	(10, 9, 7, 8)	$(35, 13)^d$	2	7	2.51	0.59	7.97	12
	(A,Ep	(6, 9, 11, 12)	(8, 46)	5	7	4.15	1.12	5.47	29
$\mathbf{A}_{0}\mathbf{E}$	EA ET	(7, 9, 11, 7)	$(14, 42)^d$	3	8	3.01	0.93	8.08	20
v	$\mathbf{E} \left\{ egin{array}{l} \mathbf{A_o^E_R} \\ \mathbf{A_o^E_L} \\ \mathbf{A_o^E_o} \end{array} \right.$	(7, 8, 9, 9)	$(13, 29)^d$	6	6	3.04	1.34	7.59	31
	$A_R A_R + A_R E_R^c$	(13, 15, 13, 16)	(62, 63)	19	261	2.68	1.04	0.86	912
			•			$(2.78, 0.93)^c$	(0.9, -)	(0.74, 3.11)	(866, 46)
	$A_R A_0$	(13, 15, 6, 8)	(62, 7)	5	28	1.43	0.62	1.10	67
	$\mathbf{A_o}\mathbf{A_R}$	(6, 8, 13, 15)	(7, 62)	5	19	1.56	0.80	1.59	54
	$A_0A_0$	(6, 8, 6, 8)	(7, 7)	2	8	2.79	0.47	0.53	15

<sup>a</sup> Number of hypercubes in peak. <sup>b</sup> Height of maximum (from Table III). <sup>c</sup> See text for discussion of these peaks, which contain two types of structure.  $^d$  Square 25 is structurally adjacent to the  $E_R$  region because of the second discontinuity mentioned in section II of the text. Squares 13 and 14 do not fall within the  $A_0$  island of Figure 2 but are structurally adjacent to it.

ferent peak types. The first peak listed, with maximum at (25, 42) (contracted coordinates), contains a number of dissimilar types of extended structure, i.e., structures with different handedness (not shown explicitly in the table), and arises because of the mixing of  $p_i p_{i+1} = +1$  and  $p_i p_{i+1}$ = -1 structures in this limited region of  $G^{(5)}$ . (For this reason, this peak is denoted only as EE in Table IV.) It will be seen from the analysis of the transformed version of G<sup>(5)</sup> (Appendix III) that the proper assignment of the contents of this peak does not affect any of the conclusions of this section. We also observe that the peak with maximum at (34, 29) contains two types of structure,  $E_0E_0$  and E<sub>L</sub>E<sub>0</sub>, with respective populations of 328 and 116. The fact that these are contained in a single peak indicates that the two structural types form a continuum, with all intermediate structural types between the two extremes ( $E_0E_0$  and  $E_{\rm I}E_0$ ) occurring in substantial numbers. This undoubtedly corresponds to the fact (noted in section II) that a similar continuum connecting  $E_L$  and  $E_0$  types occurs in  $\mathbf{G}^{(3)}$ . (The occurrence of two structural types in separate peaks indicates that structures intermediate in character between the two occur with lesser frequency than the two extreme structures.)

The populations of the structurally defined EE peaks (given in the last column of Table IV) are ordered as follows:  $(E_0E_0) > E_0E_L > E_LE_L > (E_LE_0) > E_RE_L > E_LE_R$  $= E_R E_0 > E_R E_R$ . Two types are placed in parentheses to indicate that they occur in the same peak and therefore are divided by inspection into the two classes. It is seen from the above inequality that structures which are made up of various combinations of  $E_0$  and  $E_L$  four- $C^{\alpha}$  units predominate in the extended region on the five- $C^{\alpha}$  scale. There is a strong preference for nearly flat  $E_0E_0$  structures.  $E_0E_L$  structures are markedly preferred to  $E_LE_0$  types,

which occur with nearly the same frequency as E<sub>L</sub>E<sub>L</sub> structures. Structures containing E<sub>R</sub> units are considerably less common. Within this group, the most frequently occurring structure is the E<sub>R</sub>E<sub>L</sub> type, which is heavily favored over the  $E_L E_R$  type. This latter occurs with the same frequency as the E<sub>R</sub>E<sub>0</sub> structure, and the least common feature is the  $E_R E_R$  five- $C^{\alpha}$  unit.

It is of interest to note that, in mixed EE structures, E<sub>L</sub> four- $C^{\alpha}$  units show a strong preference for the second position (E<sub>X</sub>E<sub>L</sub>), whereas E<sub>R</sub> structures prefer the first position  $(E_R E_X)$ , where X is R, L, or 0.

All of the peaks in the EE region have positive average correlation between the component four- $C^{\alpha}$  units ( $\tilde{c} > 1$ ), i.e., they pair with greater-than-random frequency, confirming the tendency of E units to associate to form extended strands. Although E<sub>R</sub> structures occur least frequently, relative to their population, they are the most likely to associate with other E<sub>X</sub> structures; E<sub>0</sub> and E<sub>L</sub> (each associated with other  $E_x$ ) structures occur even less frequently relative to their populations. The width of the peaks, as measured by  $\bar{d}$ , varies from 0 for the  $E_R E_R$  peak at (54, 46), which includes a single hypercube, to 1.60 for the  $E_0E_0 + E_LE_0$  peak at (34, 29). As indicated in Table IV, it is possible to calculate separate values of the various parameters for the two components of this peak (given in parentheses). The value of  $\bar{d}$  for the  $E_L E_0$  component is not meaningful, because the peak contains two kinds of structures, and the maximum is not contained in the  $E_L E_0$ region of the peak. The value of d for the  $E_0E_0$  component, however, is meaningful; it is 1.46, making this the broadest peak in  $G^{(5)}$ . This is followed closely by the  $E_L E_L$  peak at (41, 42) with d = 1.43. These peaks therefore contain a relatively broad range of structural types. The peaks then narrow (i.e., decrease in  $\bar{d}$ ) in the order  $E_0E_L$ ,  $E_RE_L$ ,  $E_RE_0$ ,  $E_L E_R$ ,  $E_R E_R$ . The large difference in  $\bar{d}$  between the  $E_R E_L$  and  $E_L E_R$  peaks corresponds to their difference in population.

The most frequently occurring structures, those constituting the  $E_0E_0$  component of the peak at (34, 29), show a strong tendency to be repeating, with  $\bar{\Delta} = 1.09$ . This is also true of the  $E_L E_L$  peak, with  $\bar{\Delta} = 1.57$ , and of the very small  $E_R E_R$  peak, with  $\bar{\Delta} = 1.41$ . The  $E_L E_0$  and  $E_0 E_L$ peaks, which represent transitions between twisted and flat structures, have intermediate values of  $\bar{\Delta}$  at 3.24 and 2.93, respectively. The small E<sub>R</sub>E<sub>0</sub> peak is somewhat less repeating, with  $\bar{\Delta}$  = 4.80. The  $E_R E_L$  and  $E_L E_R$  peaks, which mix twisted four-C<sup>α</sup> structures of opposite handedness, have values of  $\bar{\Delta}$  of 6.22 and 5.41, respectively. We see then that extended structures which are almost repeating predominate; this is in keeping with the prevalent intuitive picture of an extended strand. There are, however, substantial numbers of extended structures which are not repeating structures.

The next two groups of peaks in Table IV represent five- $C^{\alpha}$  structures in which the various types of extended four- $C^{\alpha}$  units combine with four- $C^{\alpha}$  segments which fall in the A<sub>R</sub> region. Peaks corresponding to all possible E<sub>X</sub>A<sub>R</sub> and A<sub>R</sub>E<sub>X</sub> structures are observed. The populations of the three  $EA_R$  peaks are ordered as  $\underline{E}_LA_R > \underline{E}_0A_R > \underline{E}_RA_R,$  with a marked preponderance of  $E_LA_R$  over the other two possibilities. The  $E_0A_R$  and  $E_LA_R$  peaks are rather narrow  $(\bar{d}=0.81,\,1.00)$ . The  $E_RA_R$  peak is less populated but of equal width  $(\bar{d}=1.03)$ . The values of  $\bar{\Delta}$  are reasonable for structures which are clearly not repeating. The most nearly repeating structure is contained in the  $E_R A_R$  peak, with  $\bar{\Delta}$  = 3.66. The most striking feature of these peaks is the fact that, for almost all of them,  $\bar{c} \lesssim 1$ , indicating that (on the five- $C^{\alpha}$  length scale) E and  $A_R$  structures tend to avoid one another. [The sole exception is the small  $A_R E_0$  peak at (55, 35), which is actually distorted from the classical  $A_R$  region toward an  $E_R E_0$  type of structure.] This is reasonable on steric grounds. One expects that the primary circumstance under which E and A (bend) structures pair should be at turns involved in  $\beta$ -sheet structures. Clearly, AR bends, which are quite nonplanar, are not favored under conditions where it is necessary to ensure the correct registration of the extended strands on either side of the bend.

The same effect is observed in the group of  $A_RE$  peaks. This includes the three peaks listed under the  $A_RE$  heading, as well as the  $A_RE_R$  component of the enormous  $A_RA_R$  peak, which is separated by inspection in the same manner as the  $E_0E_0+E_LE_0$  peak above. The values of  $\bar{c}$  of these are less than 1, with the exception of the  $A_RE_0$  peak, as noted above. It should be noted that the number of  $A_RE$  structures is considerably less than the number of  $EA_R$  structures and that the peaks are narrow.

In contrast to the above observations, the values of  $\bar{c}$  of the  $A_0E_X$  and  $E_0A_0$  peaks are all quite large. It is clear that flat bends prefer to associate with extended structures, probably because of the role that these bends play in the construction of  $\beta$  structure. The  $A_0E$  peaks are all broader  $(0.93 \leq \bar{d} \leq 1.34)$  than the narrow  $(\bar{d}=0.59)$   $E_0A_0$  peak. This correlates with the marked asymmetry in population, with  $A_0E$  structures being much more common. It is also of interest that  $E_RA_0$  and  $E_LA_0$  structures occur with very low frequency—so low, in fact, that corresponding peaks do not appear in  $\mathbf{G}^{(5)}$ .

The remaining peaks in  $G^{(5)}$  involve combinations of the two bend types. The  $A_RA_R$  peak at (62, 63) involves a minor  $A_RE_R$  component, discussed above. The enormous helical component has positive correlation ( $\bar{c}=2.78$ ), very

sharp definition ( $\bar{d} = 0.9$ ) considering its population (866), and is, of course, a repeating structure ( $\bar{\Delta} = 0.74$ ).

The two peaks involving combinations of the  $A_0$  and  $A_R$  four-C° structures show positive, but rather low, correlation. Clearly, these structures are substantially less favored than  $A_RA_R$  structures, despite the fact that they are related structurally rather closely. A measure of this relation is given by the values of  $\bar{\Delta}$  [1.10 for the  $A_RA_0$  peak at (62, 7) and 1.59 for the  $A_0A_R$  peak at (7, 62)], which indicate that these structures are substantially less repeating than  $A_RA_R$  ( $\bar{\Delta}=0.74$ ). Both peaks are fairly narrow, with  $\bar{d}=0.62$  ( $A_RA_0$ ) and 0.80 ( $A_0A_R$ ).

The last peak in  $G^{(5)}$  is the  $A_0A_0$  peak at (7, 7), which shows substantial positive correlation  $(\bar{c} = 2.79)$ . The peak is quite narrow  $(\bar{d} = 0.47)$ , and the  $\bar{\Delta}$  value of 0.53 indicates that the points in this peak represent quite strongly repeating structures.

The foregoing results indicate that there are substantial differences in the behavior of  $A_R$  and  $A_0$  four- $C^{\alpha}$  structures on the five- $C^{\alpha}$  length scale, despite the fact that they are rather closely related structurally. In fact, they seem to exhibit complementarity in their interactions with neighboring structures; in a given situation, one type of bend is preferred, and the other tends to be avoided. In combination with extended structures, AR bends exhibit negative correlation, and A<sub>0</sub> structures positive correlation. In combination with other A<sub>R</sub> structures, A<sub>R</sub> shows highly positive, and A<sub>0</sub> substantially less positive, correlation. We see, then, that the selectivity which seemed to be absent, or very weak, in the choice of possible nucleating structures on the four- $C^{\alpha}$  length scale begins to become more marked on the five- $C^{\alpha}$  scale. Consideration of the four- $C^{\alpha}$  structures contained in G(3) revealed that any conformation accessible to a four- $C^{\alpha}$  unit is similar to some structure which occurs with greater-than-average frequency and is therefore convertible with minimal deformation into a potential nucleating structure on the four- $C^{\alpha}$  scale. Consideration of  $G^{(5)}$ , however, reveals that there are restrictions on the manner in which these four- $C^{\alpha}$  nucleation structures can coalesce. Of greatest significance is the fact that two major four- $C^{\alpha}$  nucleation structures, the E and A<sub>R</sub> structures, show a pronounced tendency to avoid one another. The structure which shows a positive tendency to associate with both, the A<sub>0</sub> structure, occurs less fre-

From the foregoing considerations, a hypothesis about protein folding can be formulated. As renaturing conditions are created, nuclei of relatively long-lived structure are formed. These begin as very local, single-residue structures and lengthen as conditions become more renaturing. These structures are likely to be those which are most strongly stabilized by local interactions and therefore also appear with highest frequency in the native structure.

Our analysis suggests that the nucleating structures which form up to the five- $C^{\alpha}$  length scale fall into either of two classes: (a) essentially repeating structures, which can be either extended strands or right-handed helices, or (b) nonrepeating structures. (As pointed out above, extended strands are not strictly repeating structures, being composed of  $E_L$ ,  $E_R$ , and  $E_0$  four- $C^{\alpha}$  units.) This is a more general result than is suggested by some previous work, 9 in which the nuclei are assumed to fall into one or the other class. Nonrepeating structures are likely to involve  $A_0$  (flat) bends in combination with either  $E_X$  or  $A_R$  structures. Nucleating structures which are combinations of  $A_R$  and  $E_X$  four- $C^{\alpha}$  structures are not likely to occur.

We have said nothing about the order in which various nuclei form in the chain, and indeed our analysis does not

Table V
Peaks in Transformed G<sup>(5)</sup>

peak type	coordinates	contracted coordinates	$N_{\mathbf{c}}$	h	$\overline{c}$	$\overline{d}$	$\overline{\Delta}$	popu- lation
/E <sub>R</sub> E <sub>L</sub>	(11, 12, 11, 7)	(46, 42)	6	9	3.05	1.18	6.24	39.5
$\mathbf{E_{L}E_{R}}$	(12, 4, 11, 12)	(50, 46)	1	5	4.23	0.0	8.06	5
$\mathbf{E_L}\mathbf{E_R}$	(11, 7, 11, 12)	(42, 46)	2	8	2.13	0.54	5.42	13
$_{\mathbf{E}_{\mathbf{E}}})\mathbf{E}_{\mathbf{L}}\mathbf{E}_{\mathbf{L}}$	(11, 6, 11, 7)	(41, 42)	19	16	2.03	1.49	1.77	147
$^{22}$ ) $E_RE_R$	(12, 13, 11, 12)	(54, 46)	1	6	3.25	0.0	1.41	6
$egin{array}{c} \mathbf{E}_{o}^{\dagger}\mathbf{E}_{L}^{\dagger} \\ \mathbf{E}_{o}\mathbf{E}_{o}^{\dagger} + \mathbf{E}_{L}\mathbf{E}_{o} \end{array}$	(9, 9, 11, 6)	(29, 41)	25	20	2.70	1.65	3.34	210
$\mathbf{E}_{0}\mathbf{E}_{0}^{-}+\mathbf{E}_{L}\mathbf{E}_{0}$	(10, 8, 9, 9)	(34, 29)	37	31	2.32	1.60	1.66	445
$(\mathbf{E}_{\mathbf{R}}^{\circ}\mathbf{E}_{\mathbf{o}}^{\circ})$	(12, 13, 9, 8)	(54, 28)	2	8	4.40	0.77	4.80	13
$(\mathbf{E}_{0}\mathbf{A}_{\mathbf{R}} + \mathbf{E}_{0}\mathbf{A}_{0})$	(10, 9, 13, 15)	(35, 62)	13	22	1.39	1.04	6.61	94
$EALA_R$	(11, 7, 13, 15)	(42, 62)	12	26	0.96	1.16	9.0	111
$(E_R A_R + E_R A_o)$	(11, 12, 13, 16)	(46, 63)	6	10	1.03	0.99	3.88	41
$-(A_R E_0)$	(12, 14, 10, 9)	(55, 35)	1	8	1.45	0.0	5.39	8 7
$A_R E \langle A_R E_R$	(13, 15, 8, 12)	(62, 25)	1	7	1.32	0.0	5.83	
$(\mathbf{A_R^E}\mathbf{E_L^F} + \mathbf{A_R^E}_0 + \mathbf{A_0^E}_L + \mathbf{A_0^E}_0)$	(13, 16, 11, 7)	(63, 42)	12	13	1.44	1.69	8.57	80.5
$A_{R}A_{R} + A_{R}E_{R} + A_{R}A_{o} + A_{o}A_{o} + A_{o}A_{R}$	(13, 15, 13, 16)	(62, 63)	29	286	2.36	1.15	1.0	1111

provide any information on this point. It should be noted, however, that none of the structures which we have proposed as five- $C^{\alpha}$  nucleating structures can be ruled out as the initial nucleating structure. At first glance it may seem unlikely that the formation of an extended strand, for example, should be a necessary first step in the refolding process. It is conceivable, however, that the separation of two parts of the molecule, which can be accomplished through the formation of a linear structure, is a necessary preliminary to correct folding in some proteins. In other cases, the formation of one or several bends may be the initial step in refolding. We thus take a somewhat more general view of nucleation than did Matheson and Scheraga, who postulate hairpin bend formation as the initial step in nucleation. This difference may be more semantic than real, however, since those authors based their work on hydrophobic interactions and therefore defined the primary (initial) nucleating structure as one which is stable and promotes a more compact chain structure, with strong hydrophobic bonding, irrespective of any repeating structure which may be present. Here, we take the view that a decrease in chain dimensions is not necessarily the first step in correct folding and therefore permit the consideration of repeating structures as initial nuclei. The (sometimes) larger structures of ref 8 form after those on the five- $C^{\alpha}$  length scale.

Similarly, our views are more general than those of Tanaka and Scheraga, 25 who postulate that the initial folding nuclei are those (including  $\alpha$  helices<sup>9</sup>) which appear near the diagonal in the protein contact map. It should be emphasized again that none of these viewpoints are excluded by our analysis. We maintain that the kinetically visible nucleating step, for example, leads to a properly folded native conformation only because of interactions brought about by other locally stabilized structures present in the chain. We therefore regard these structures as nuclei essential for proper folding. In this sense, the differential geometric analysis provides a unifying framework for the various proposed nucleation mechanisms. It also illustrates that selectivity of nucleation types increases as nucleation proceeds to longer length scales. Further, it is able to identify precisely the types of nonregular structure<sup>8,25,26</sup> that are likely to take part in nucleation, and which nonregular structures are apparently not sufficiently stabilized by purely internal interactions, and are therefore stabilized by interactions on a larger length scale (e.g., those discussed in ref 8) than that which we have explored thus far.

An interesting question raised by these results is whether there is an upper limit to the length scale on which characteristic structures can be observed which are common to a wide assortment of proteins. Intuitively one expects that such a limit does exist. Inspection of structure on a scale above this limit should reveal a broad distribution of structural types which reflect the individual folding patterns of the various proteins (or classes of proteins) in the sample. We have no quantitative answer to this question at present.

Another question of interest is the influence of residue type on local chain structure as revealed by the differential geometric representation. The results presented above are averaged over the amino acid sequence; one expects that there will be variations as a function of amino acid composition. This problem will be addressed in a forthcoming publication.

## IV. Summary

We have used a graphical method based on the differential-geometric approach to investigate the backbone structure of proteins on the four- and five- $C^{\alpha}$  length scales. The result of this investigation led to a number of conclusions about protein structure and protein folding.

- 1. On the four- $C^\alpha$  length scale, certain distinct structural types are observed. The extended region shows two peaks, one corresponding to right-handed twisted four- $C^\alpha$  structures  $(E_R)$  and one containing left-handed twisted  $(E_L)$  structures and nearly flat  $(E_0)$  structures. There is a structural continuum embracing the latter two, which is separate from the smaller  $E_R$  peak. Another structural continuum observed on the four- $C^\alpha$  scale involves the very common  $\alpha$ -helical  $(A_R)$  and less common nearly flat  $(A_0)$  bend types.
- 2. An analysis of the types of structure which occur with high frequency reveals that such structures are distributed throughout the occupied region on the four- $C^{\alpha}$  length scale. Therefore, any structure accessible to a four- $C^{\alpha}$  unit is close to some such structure, probably within fluctuation range. It seems reasonable that such high-frequency structures are energetically favored and are therefore potential nucleation structures in the denatured protein. Thus, potential nucleation structures are distributed throughout the accessible region on the four- $C^{\alpha}$  length scale, and there is rather weak selectivity in nucleation on this length scale.
- 3. Extension of this analysis to five- $C^{\alpha}$  structures reveals distinct peaks corresponding to various combinations of the four- $C^{\alpha}$  structures noted in (1) above. The extended region is particularly rich in peak structure, with peaks corresponding to most of the possible combinations of  $E_L$ ,

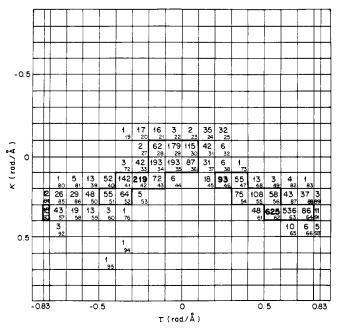


Figure 3. Transformed single-site distribution of values of  $(\kappa, \tau)$  for the same protein sample. Notations as in Figure 2. It should be noted that the lines  $\tau = \pm 0.83$  are equivalent because of a cyclic boundary condition. The numbering of the squares (small numbers) is the same as in Figure 2; squares with numbers greater than 76 are not occupied in the untransformed distribution of Figure 2.

 $E_R$ , and  $E_0$  structures. There are wide variations in the frequency with which these structures occur. Other peaks are observed corresponding to combinations of extended structures with  $A_0$  and  $A_R$  bend types, to  $A_R$ -helical segments, to combinations of  $A_0$  and  $A_R$  bends and to  $A_0A_0$  structures. Analysis of the correlation between neighboring structures reveals that there is a tendency for certain structures to avoid association and for others to associate preferentially. This leads to complementarity between similar structural types. For example, the two bend types are complementary in their association with extended structures;  $A_0$  bends tend to associate with extended four- $C^\alpha$  structures and  $A_R$  bends to avoid them.

4. We suggest that those five- $C^{\alpha}$  structures which exhibit positive correlation between their component four- $C^{\alpha}$ structures are potential five- $C^{\alpha}$  components of nucleating structures. These are the E<sub>X</sub>E<sub>Y</sub>, E<sub>0</sub>A<sub>0</sub>, A<sub>0</sub>E, A<sub>R</sub>A<sub>R</sub>, and A<sub>0</sub>A<sub>0</sub> structures. A<sub>R</sub>A<sub>0</sub> and A<sub>0</sub>A<sub>R</sub> show positive, but considerably lower, correlation. Clearly there is a higher level of selectivity in nucleation on the five- $C^{\alpha}$  length scale, brought about by the preferential association of certain structures. There is also a considerable difference in the roles played by the various potential nucleation structures. The EE and  $A_RA_R$  structures seem to provide nuclei for the formation of repeating extended or helical structures, respectively. The role of the A<sub>0</sub> structure is important despite its numerical inferiority, since it seems to provide the principal building block (in combination with either E or A<sub>R</sub> units) for nucleation of nonregular backbone structures, which are crucial for the formation of the final globular conformation of the protein molecule.

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# Appendix I. Analysis of Transformed Single-Site Distribution

In Figure 3, we show the transformed version of Figure 2; i.e., all points of Figure 2 with  $p_i p_{i+1} = +1$  have been

transferred (in Figure 3) to squares with  $p_i p_{i+1} = -1$  by the methods outlined in paper 2.5 An analysis exactly analogous to that performed for Figure 2 in the main text gives two islands of greater-than-random occupancy, which are indicated in the figure by heavy outlines. It can be seen (by comparing Figures 2 and 3) that the island including the E<sub>L</sub> and E<sub>0</sub> structures remains unaltered. The island containing the E<sub>R</sub> and A<sub>R</sub> peaks is now expanded, due to the transposition of the A<sub>0</sub> region to positions adjacent to the A<sub>R</sub> and A<sub>L</sub> regions. It should be noted that square 64, which was not part of this island in Figure 2, is now included (since it now contains points which fell in the A<sub>0</sub> region in Figure 2) and the population of square 63 is considerably increased. It should also be noted that the partial square 90 is also a region of marginally high frequency, corresponding to square 9 of Figure 2. It can thus be seen that the same structures which appear in the high-frequency islands of Figure 2 appear in those of Figure 3. Figure 3 also brings out the close structural relationship between the A<sub>R</sub>, A<sub>0</sub>, and A<sub>L</sub> regions. (It should be noted that in Figure 3 there is a periodic boundary condition operating between the lines  $\tau = \pm 0.83$ , so that, for example, "squares" 90 and 91 are adjacent.)

# Appendix II. Procedure for Separating Islands into Peaks

The following procedure is valid in any dimension and reduces in the three-dimensional case to an intuitively clear procedure.

(1) For each maximum on the island, all squares are identified which are accessible from the maximum by proceeding downward or horizontally (in population). In general, there will be squares which are accessible from more than one peak. Therefore, (2) each square which is accessible from more than one peak is assigned to the peak to which it is closest; except that (3) squares which are nearest neighbors of two peaks simultaneously (in a sense to be defined) are divided equally between the two peaks.

The sense in which we define nearest neighbors is that two squares, for example,  $(m_1, n_1)$  and  $(m_2, n_2)$  in a three-dimensional graph, are nearest neighbors if

$$\max (|m_1 - m_2|, |n_1 - n_2|) = 1$$

As pointed out in section III of the main text, each square in  $\mathbf{P}^{(2)}$  has 8 nearest neighbors in this sense, whereas each hypercube in  $\mathbf{P}^{(4)}$  has 80 nearest neighbors.

# Appendix III. Analysis of Transformed Version of $G^{(5)}$

Table V shows the peaks in the transformed version of  $G^{(5)}$  organized by type, in analogy with Table IV. It can be seen that the essential change in the extended region brought about by the transformation is the redistribution of the contents of the mixed peak [at (25, 42) of Table IV], principally to the  $E_LE_L$  peak at (41, 42) and the  $E_0E_L$  peak at (29, 41). This increases the values of  $\bar{d}$  and  $\bar{\Delta}$  of these peaks somewhat. The conclusions drawn in the main text remain unaffected.

The principal effect of the transformation is to place the  $A_0$  region next to the  $A_R$  and  $A_L$  regions of the  $(\kappa, \tau)$  plane. This results in the elimination from Table V of separate peaks containing the  $A_0$  structures. The peaks centered at the  $A_R$  region are broader and contain  $A_0$  structures as well as  $A_R$  structures. The values of  $\bar{c}$  of these peaks are intermediate between those of the corresponding separate peaks in Table IV, reflecting the complementarity between  $A_R$  and  $A_0$  structures noted in the main text (section III).

It is clear that the use of the untransformed version of  $G^{(5)}$  enables us to separate the roles of  $A_R$  and  $A_0$  structures

conveniently, with no loss of information.

#### References and Notes

- (1) This work was supported by research grants from the National Science Foundation (Grant No. PCM79-20279) and from the National Institute of General Medical Sciences of the National Institutes of Health, U.S. Public Health Service (Grant No. GM-14312).
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  (11) By definition, 5,10 p<sub>i</sub>p<sub>i+1</sub> = +1 for A<sub>0</sub> bends.
  (12) In order to assess the significance of the observed difference between the three- and five-dimensional distributions (of parts II and III, respectively) and the uniform distributions to which they are compared in the text,  $\chi^2$  goodness-of-fit tests were carried out. In both cases it was determined that the distributions differ from the corresponding uniform distribution with greater than 0.999 significance.
- (13) It should be noted that the discussion of the parity (i.e., handedness) of the extended region in paper 1 is in error, due to the error in the assignment of parity regions in that paper; a corrected presentation is given in paper 2. In fact, it is known<sup>14,15</sup> that, while  $\beta$  sheets have predominantly righthanded twists, the extended strands from which they are constituted are mainly left-handed in their twist. Our observations in this paper are in agreement with this earlier conclusion. 14,15
- (14) Chothia, C. J. Mol. Biol. 1973, 75, 295.
- (15) There has been some ambiguity in the literature concerning the definition of the handedness of individual strands in  $\beta$ sheets. We define handedness in a manner consistent with the nomenclature for the conformations of polypeptide chains;16 i.e., the sign of the dihedral angle around a virtual bond is defined in the same way as that around a chemical bond in Figure 1 of ref 16. With this definition, all of the individual strands in the  $\beta$  sheets commonly observed in proteins have a predominantly left-handed twist. Our definition of parity is in agreement with this definition of handedness. Also, this definition of handedness is consistent with that derived earlier<sup>17</sup> from an analysis of curves of constant n (where n is the number of residues per turn of helical structures) on a  $(\phi, \psi)$ diagram: the dihedral angles of the individual strands in the  $\beta$  sheets observed in proteins fall predominantly in the region of the  $(\phi, \psi)$  map that corresponds to left-handedness. 14,18 Furthermore, if one looks at successive backbone oxygen atoms of these individual strands, the line connecting them forms a left-handed helix around the backbone. If one considers every second backbone oxygen atom, however, then these form a right-handed helix around the backbone. The direction of hydrogen bonding between two adjacent strands in a  $\beta$  sheet is determined primarily by the directions of every second C=O

bond of each strand, and hence the curved sheet passing through the successive hydrogen bonds between two strands in the commonly observed  $\beta$  structures in proteins has a right-handed twist. In other words, left-handed  $\beta$  strands combine to form a  $\beta$  sheet with a right-handed twist, when viewed along the strands. This relationship between the twist of the strands and of the sheet has been demonstrated in Figure 4 of ref 18, which used the same convention as that adopted here. An analogous relationship is seen in collagen and collagen-like triple-stranded polypeptides: the individual strands form left-handed helices, but the triple-stranded structure has a right-handed twist around its major axis. 19 In several other studies,14 the orientation of every second residue, rather than of every residue, has been used to define the twist of the individual strands, leading to a description opposite to that used here. As a separate but related point, if a curved sheet has a right-handed twist (as described at the beginning of this footnote) when viewed along the direction of the strands, then the twist of the sheet is opposite, i.e., left-handed, when one looks at the sheet perpendicularly to the direction of the chains (see Figure 4 of ref 18).

(16) IUPAC-IUB Commission on Biochemical Nomenclature Bio-

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- (20) It is important to remark, in connection with eq 1, that the graphic approach introduces a coarse graining into the distribution and therefore into any average quantity  $\bar{D}$  calculated from eq 1; of course, such coarse graining is not present in the actual distribution. For example, all points in the  $(\kappa, \tau)$  plane which fall in a given square of the graph will be considered to be located at the center of that square for the purpose of calculating distances. Thus, two points which are very close together, but which fall in adjacent squares of the graph, are regarded as being 0.1 rad/Å apart. The use of a finer grid for the graph can reduce the effect of this approximation, but a compromise must be struck between the need to sharpen the graph and the inherent limit in the distinctness of points due to experimental error in the X-ray coordinates from which the distribution is constructed. It seems that the  $(0.1 \times 0.1)$  rad/Å
- grid represents a good compromise between these two effects.

  (21) Equation 2 is adequate for measuring distances within localized structures such as peaks in Figure 2. In general, the two discontinuities discussed in the text must be taken into account when calculating distances between two points in the  $(\kappa,$ τ) plane, and this is done where appropriate throughout this work.
- (22) This is in contrast to previous work<sup>23,24</sup> in which we excluded  $\alpha_R$  helices from the structures identified as bends. The present work demonstrates the structural identity which actually obtains.
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- There are only 3729 nearest-neighbor pairs in the 22 proteins because ribonuclease S has one break in the backbone and  $\alpha$ -chymotrypsin has two breaks.